

## WS 4.2 Meander line

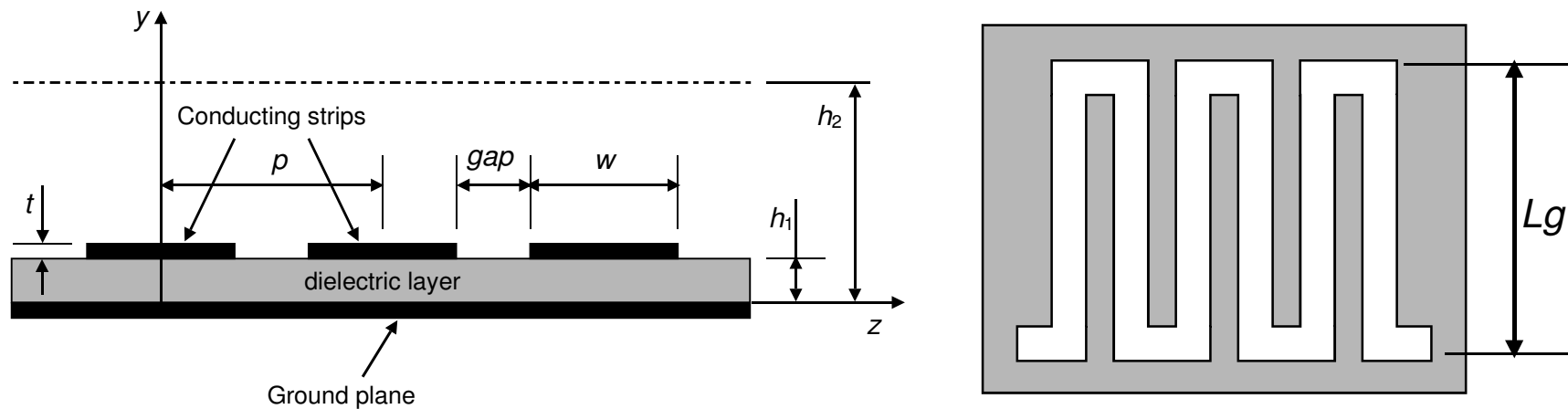
© 2017 Richard G Carter

This Mathcad 14 worksheet is designed to accompany the author's book "Microwave and RF Vacuum Electronic Power Sources", Cambridge University Press (2018). The section, equation, and figure numbers refer to the corresponding sections, equations, and figures in the book. Data input fields are highlighted in yellow and output fields are highlighted in green.

This resource is provided free of charge by Cambridge University Press with permission of the author, but is subject to copyright. You are permitted to view, print and download this resource for your own personal use only, provided any copyright lines are not removed or altered in any way. Any other use, including but not limited to, distribution of the resource in modified form, or via electronic or other media, is strictly prohibited unless you have permission from the author and provided you give appropriate acknowledgement of the source.

The contents of this sheet are provided for educational purposes only and no warranty is expressed or implied that they are suitable for use as professional design tools.

J. A. Weiss, "Dispersion and Field Analysis of a Microstrip Meander-Line Slow-Wave Structure," *Microwave Theory and Techniques, IEEE Transactions on*, vol. 22, pp. 1194-1201, 1974. Note: here  $\phi$  is the phase change between adjacent lines and not the phase change per pitch as used by Weiss. The dimensions in W. E. Courtney, "Printed-circuit RF-keyed crossed-field amplifier," Lincoln Laboratory, Lexington, MA. 1975 are apparently the same.



**Structure dimensions**

$h_1 := 0.062 \cdot \text{in}$

$h_2 := 2.6 \cdot h_1$

$L_g := 0.537 \cdot \text{in}$

$w := 0.045 \cdot \text{in}$

$\text{gap} := 0.057 \cdot \text{in}$

$p := w + \text{gap}$

$\epsilon_2 := 6.5$

$p = 2.591 \cdot \text{mm}$

$h_1 = 1.57 \cdot \text{mm}$

$L_g = 13.6 \cdot \text{mm}$

$w = 1.14 \cdot \text{mm}$

$\text{gap} = 1.45 \cdot \text{mm}$

$p = 2.59 \cdot \text{mm}$

$h_2 = 4.09 \cdot \text{mm}$

Choose mesh size for FD calculations

$\text{mesh} := 0.142 \cdot \text{mm}$

Note: In this model  $t = 0$ 

$NH1 := \frac{h_1}{\text{mesh}}$

$NH2 := \frac{h_2}{\text{mesh}}$

$NP := \frac{p}{\text{mesh}}$

$NW := \frac{w}{\text{mesh}}$

$NH1 = 11.09$

$NH2 = 28.834$

$NP = 18.245$

$NW = 8.049$

Note: these numbers must be close to integers

$Nh1 := \text{round}(NH1)$

$Nh2 := \text{round}(NH2)$

$Np := \text{round}(NP)$

$Nw := \text{round}(NW)$

$Nw2 := \frac{Nw}{2}$

$Np2 := \frac{Np}{2}$

$Nh1 = 11$

$Nh2 = 29$

$Np = 18$

$Nw = 8$

NB:  $Np$  and  $Nw$  must be even numbers

This variable is 0 for a conducting top plane and 2 for a symmetry boundary

$H2_{\text{sym}} := 2$

Set maximum number of iterations

$Niter := 1000$

**Calculation of capacitances using a finite difference solution of Laplace's equation**

```

CM( $\Phi, \varepsilon_2$ ) :=
  for i  $\in$  0..Np
    for j  $\in$  0..Nh2
       $V_{i,j} \leftarrow 0$ 
    V1  $\leftarrow$  100
     $V_2 \leftarrow \begin{cases} 100 & \text{if } \Phi = 0 \\ 0 & \text{if } \Phi = 0.5 \cdot \pi \\ (-100) & \text{if } \Phi = \pi \end{cases}$ 
     $R_{\text{sym}} \leftarrow \begin{cases} 2 & \text{if } \Phi = 0 \\ 0 & \text{if } \Phi = 0.5 \cdot \pi \\ 2 & \text{if } \Phi = \pi \end{cases}$ 
    for n  $\in$  1..Niter
      for i  $\in$  0..Np
         $V_{i,0} \leftarrow 0$ 
        for j  $\in$  1..(Nh1 - 1)
           $V_{0,j} \leftarrow (2 \cdot V_{1,j} + V_{0,j+1} + V_{0,j-1}) \cdot 0.25$ 
           $V_{Np,j} \leftarrow (R_{\text{sym}} \cdot V_{Np-1,j} + V_{Np,j+1} + V_{Np,j-1}) \cdot 0.25$ 
          for i  $\in$  1..Np - 1
             $V_{i,j} \leftarrow (V_{i+1,j} + V_{i,j+1} + V_{i-1,j} + V_{i,j-1}) \cdot 0.25$ 
        for i  $\in$  0..Nw2
           $V_{i,Nh1} \leftarrow V1$ 
        for i  $\in$  (Nw2 + 1)..(Np - Nw2 - 1)
           $V_{i,Nh1} \leftarrow \frac{V_{i,Nh1-1} \cdot \varepsilon_2 + 0.5 \cdot (V_{i-1,Nh1} + V_{i+1,Nh1}) \cdot (\varepsilon_2 + 1) + V_{i,Nh1+1}}{2 \cdot \varepsilon_2 + 2}$ 

```

```

for i ∈ (Np - Nw2) .. Np
  Vi,Nh1 ← V2
for j ∈ (Nh1 + 1) .. (Nh2 - 1)
  V0,j ← (2·V1,j + V0,j+1 + V0,j-1)·0.25
  VNp,j ← (Rsym·VNp-1,j + VNp,j+1 + VNp,j-1)·0.25
  for i ∈ 1 .. Np - 1
    Vi,j ← (Vi+1,j + Vi,j+1 + Vi-1,j + Vi,j-1)·0.25
V0,Nh2 ← (2·V1,Nh2 + H2sym·V0,Nh2-1)·0.25
VNp,Nh2 ← (Rsym·VNp-1,Nh2 + H2sym·VNp,Nh2-1)·0.25
for i ∈ 1 .. Np - 1
  Vi,Nh2 ← (Vi+1,Nh2 + Vi-1,Nh2 + H2sym·Vi,Nh2-1)·0.25
Q1 ←  $\left[ \sum_{i=0}^{0.5 \cdot Np} V_{i,1} - \left( \frac{V_{0,1} + V_{Np2,1}}{2} \right) \right] \cdot \varepsilon_2$ 
Q2 ←  $\left[ \sum_{j=0}^{Nh1} (V_{Np2-1,j} - V_{Np2+1,j}) - \frac{(V_{Np2-1,0} - V_{Np2+1,0})}{2} - \left( \frac{V_{Np2-1,Nh1} - V_{Np2+1,Nh1}}{2} \right) \right] \cdot \frac{\varepsilon_2}{2}$ 
Q3 ←  $\left[ \sum_{j=Nh1}^{Nh2} (V_{Np2-1,j} - V_{Np2+1,j}) - \frac{(V_{Np2-1,Nh1} - V_{Np2+1,Nh1})}{2} - \left( \frac{V_{Np2-1,Nh2} - V_{Np2+1,Nh2}}{2} \right) \right] \cdot \frac{1}{2}$ 
Q4 ←  $\left[ \sum_{i=0}^{0.5 \cdot Np} V_{i,Nh2-1} - \left( \frac{V_{0,Nh2-1} + V_{Np2,Nh2-1}}{2} \right) \right] - \left[ \sum_{i=0}^{0.5 \cdot Np} V_{i,Nh2} - \left( \frac{V_{0,Nh2} + V_{Np2,Nh2}}{2} \right) \right]$ 
Cn ←  $\frac{(Q1 + Q2 + Q3 + Q4)}{50}$ 
error("Convergence failure") if n = Niter

```

```

    (break) if  $|C_n - C_{n-1}| < 0.0001$ 
  return  $C_n$ 

```

The capacitances per unit length, divided by  $\epsilon_0$  are

$$CM(0, 1) = 1.46$$

$$CM\left(\frac{\pi}{2}, 1\right) = 3.04$$

$$CM(\pi, 1) = 3.99$$

without dielectric

$$CM(0, \epsilon_2) = 9.08$$

$$CM\left(\frac{\pi}{2}, \epsilon_2\right) = 12.45$$

$$CM(\pi, \epsilon_2) = 15.18$$

with dielectric

*Free space equivalent circuit capacitances* calculated using finite differences.

The capacitances are for phase shifts 0,  $\pi/2$  and  $\pi$  per strip (and double that for the unit cell).

$$C_e := CM(0, 1) \cdot \epsilon_0$$

$$C_{\pi/2} := CM\left(\frac{\pi}{2}, 1\right) \cdot \epsilon_0$$

$$C_o := CM(\pi, 1) \cdot \epsilon_0$$

without dielectric

$$CK_e := CM(0, \epsilon_2) \cdot \epsilon_0$$

$$CK_{\pi/2} := CM\left(\frac{\pi}{2}, \epsilon_2\right) \cdot \epsilon_0$$

$$CK_o := CM(\pi, \epsilon_2) \cdot \epsilon_0$$

with dielectric



**Equivalent circuit parameters without dielectric**

$$C_0 := C_e \quad C_1 := 0.25 \cdot (C_o - C_e) \quad C_2 := 0.25 \cdot (C_{\pi 2} - C_e - 2 \cdot C_1)$$

$$C1(\phi) := C_0 + 2 \cdot C_1 \cdot (1 - \cos(\phi)) + 2 \cdot C_2 \cdot (1 - \cos(2\phi))$$

Equation 4.39

$$L1(\phi) := \frac{1}{c^2 \cdot C1(\phi)}$$

Equation 4.43

**Equivalent circuit capacitance with the dielectric substrate**

$$CK_0 := CK_e \quad CK_1 := 0.25 \cdot (CK_o - CK_e) \quad CK_2 := 0.25 \cdot (CK_{\pi 2} - CK_e - 2 \cdot CK_1)$$

$$CK(\phi) := CK_0 + 2 \cdot CK_1 \cdot (1 - \cos(\phi)) + 2 \cdot CK_2 \cdot (1 - \cos(2\phi))$$

Equation 4.39

**Effective relative permittivity**

$$\epsilon_{\text{eff}}(\phi) := \frac{CK(\phi)}{C1(\phi)}$$

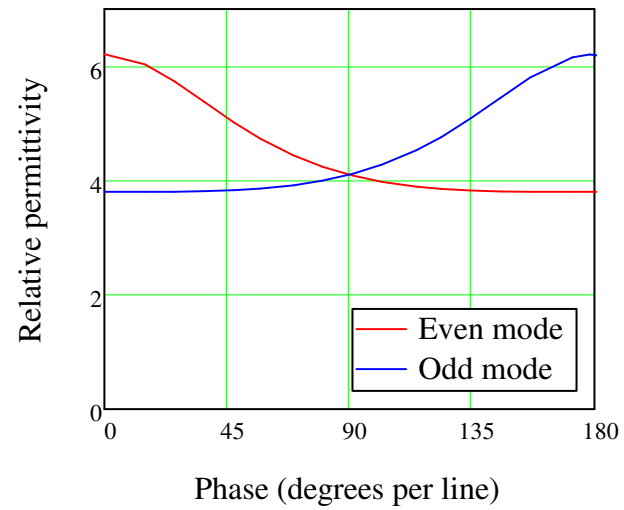
Equation 4.47

*Characteristic impedances of even and odd modes*

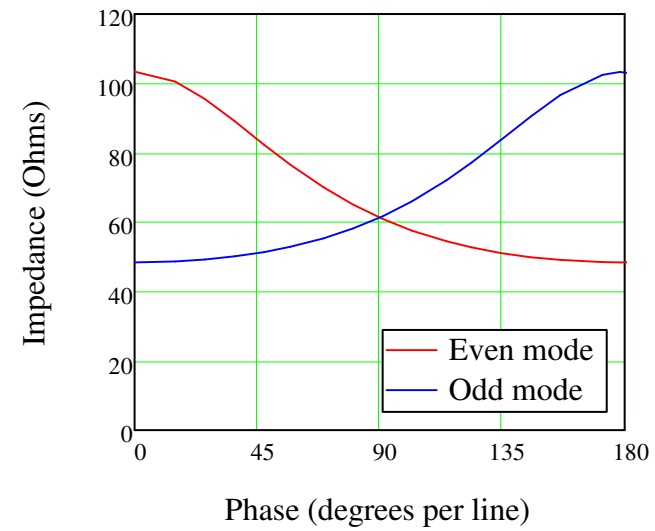
$$Z_e(\phi) := \sqrt{\frac{L1(\phi)}{CK(\phi)}} \quad Z_o(\phi) := Z_e(\pi + \phi)$$

Equation 4.50

The best agreement with Weiss is obtained when the capacitances are calculated assuming that the cover plate is a symmetry boundary



Compare Figure 5 in Weiss



Compare Figure 6 in Weiss

**Calculate the dispersion diagram**

(a) Ignoring coupling between lines and assume that the effective permittivity is that of the substrate

$$\omega_0(\phi) := \frac{c \cdot \phi}{\sqrt{(1 + \varepsilon_2) \cdot 0.5 \cdot (p + Lg)}} \quad f_0(\phi) := \frac{\omega_0(\phi)}{2 \cdot \pi \cdot \text{GHz}} \quad f_0(\pi) = 4.769 \quad \text{Path length is } L + p$$

(b) Including coupling between lines the phase velocities for the even and odd modes are

$$v_{pe}(\phi) := \frac{c}{\sqrt{\varepsilon_{eff}(\phi)}} \quad v_{po}(\phi) := \frac{c}{\sqrt{\varepsilon_{eff}(\pi + \phi)}} \quad \boxed{\text{Equation 4.48}}$$

Define the functions in Equation 4.57 for the lower and upper branches of the dispersion diagram

$$F1(\omega, \phi) := \tan\left(\frac{\phi}{2}\right)^2 - \frac{Z_o(\phi)}{Z_e(\phi)} \cdot \tan\left(\frac{\omega \cdot Lg}{2 \cdot v_{pe}(\phi)}\right) \cdot \tan\left(\frac{\omega \cdot Lg}{2 \cdot v_{po}(\phi)}\right)$$

$$F2(\omega, \phi) := \tan\left(\frac{\phi}{2}\right)^2 - \frac{Z_o(\phi)}{Z_e(\phi)} \cdot \cot\left(\frac{\omega \cdot Lg}{2 \cdot v_{pe}(\phi)}\right) \cdot \cot\left(\frac{\omega \cdot Lg}{2 \cdot v_{po}(\phi)}\right)$$

$$F3(\omega, \phi) := \cot\left(\frac{\phi}{2}\right)^2 - \frac{Z_e(\phi)}{Z_o(\phi)} \cdot \cot\left(\frac{\omega \cdot Lg}{2 \cdot v_{pe}(\phi)}\right) \cdot \cot\left(\frac{\omega \cdot Lg}{2 \cdot v_{po}(\phi)}\right)$$

$$F4(\omega, \phi) := \cot\left(\frac{\phi}{2}\right)^2 - \frac{Z_e(\phi)}{Z_o(\phi)} \cdot \tan\left(\frac{\omega \cdot Lg}{2 \cdot v_{pe}(\phi)}\right) \cdot \tan\left(\frac{\omega \cdot Lg}{2 \cdot v_{po}(\phi)}\right)$$

Solve for the lower and upper branches

$$\omega_1 := 0.5 \cdot \omega_0(\pi)$$

$$\omega_2 := 1.4 \cdot \omega_0(\pi)$$

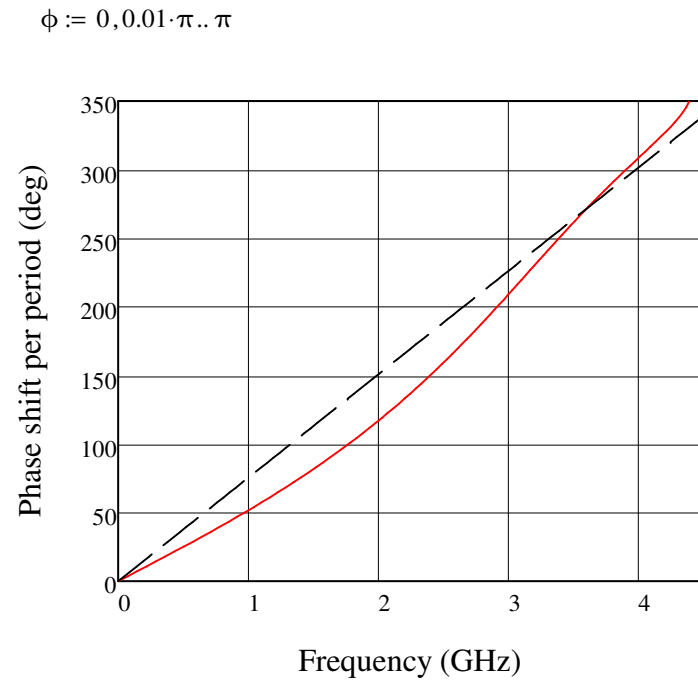
$$\omega_1(\phi) := \begin{cases} \text{root}(F1(\omega_1, \phi), \omega_1) & \text{if } \phi \leq 0.5 \cdot \pi \\ \text{root}(F3(\omega_1, \phi), \omega_1) & \text{otherwise} \end{cases}$$

$$\omega_2(\phi) := \begin{cases} \text{root}(F2(\omega_2, \phi), \omega_2) & \text{if } \phi \leq 0.5 \cdot \pi \\ \text{root}(F4(\omega_2, \phi), \omega_2) & \text{otherwise} \end{cases}$$

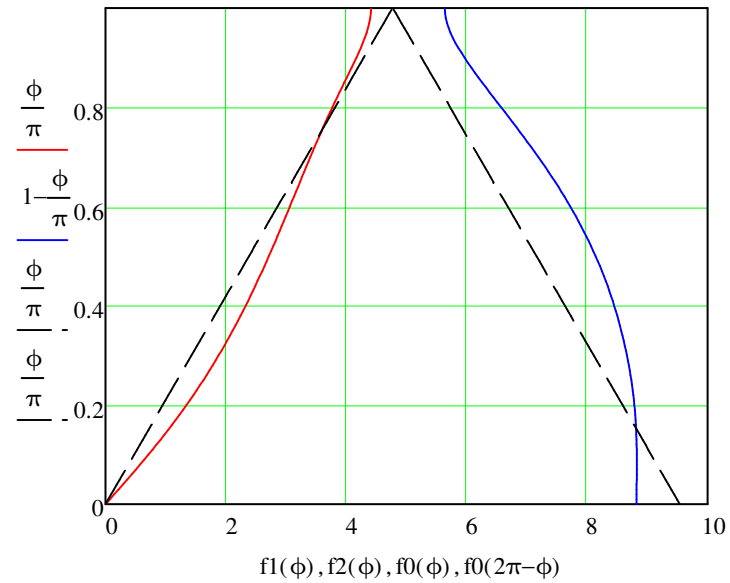
$$f1(\phi) := \frac{\omega_1(\phi)}{2 \cdot \pi \cdot \text{GHz}}$$

$$f2(\phi) := \frac{\omega_2(\phi)}{2 \cdot \pi \cdot \text{GHz}}$$





Compare Figure 6 in Courtney



Compare Figure 7 in Weiss

The best agreement with Weiss is obtained when it is assumed that the effective path length is  $L + p$  whereas he apparently chose a path length of  $L$ .  
 Note: identical results are obtained using Equation (22) in Weiss.

Calculate the image (i.e. characteristic or iterative) impedances of the lower and upper modes. (Equation 26 in Weiss).

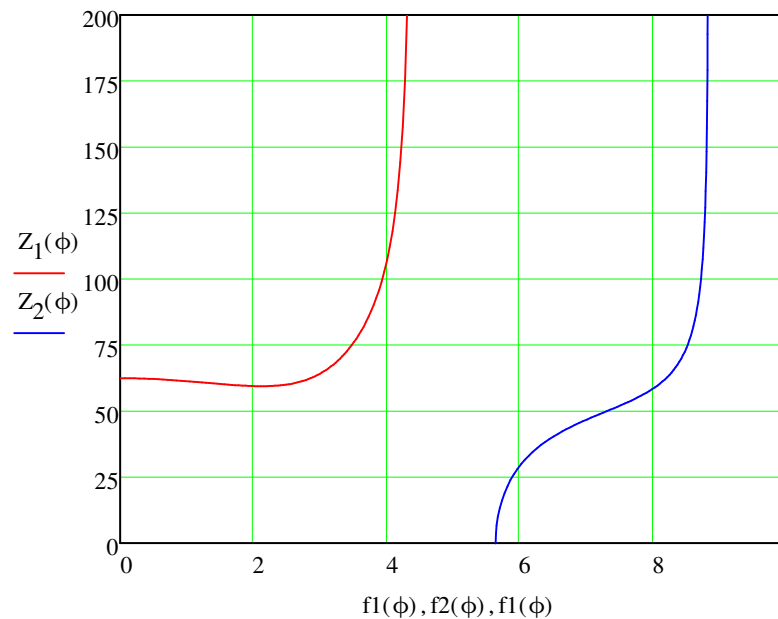
$$k_e(\phi) := \frac{\omega 1(\phi)}{v_{pe}(\phi)}$$

$$k_o(\phi) := \frac{\omega 1(\phi)}{v_{po}(\phi)}$$

For the lower band

$$Z_1(\phi) := \sqrt{Z_e(\phi) \cdot Z_o(\phi) \cdot \frac{\tan\left(\frac{k_o(\phi) \cdot L_g}{2}\right)}{\tan\left(\frac{k_e(\phi) \cdot L_g}{2}\right)}}$$

$$Z_2(\phi) := \sqrt{Z_e(\phi) \cdot Z_o(\phi) \cdot \frac{\tan\left(\frac{\omega 2(\phi) \cdot L_g}{2 \cdot v_{pe}(\phi)}\right)}{\tan\left(\frac{\omega 2(\phi) \cdot L_g}{2 \cdot v_{po}(\phi)}\right)}}$$



Compare Figure 8 in Weiss

**Calculate the Pierce impedance at height d assuming that the field is constant in the gap.**

Define the amplitude of the even mode voltage relative to ground (a in Weiss)  $V_a := 1 \cdot V$

$$b_{a1}(\phi) := -\sqrt{\frac{Z_o(\phi) \cdot \sin(k_e(\phi) \cdot Lg)}{Z_e(\phi) \cdot \sin(k_o(\phi) \cdot Lg)}} \quad \text{Weiss: equation (25)}$$

Phase velocity and group velocity  $v_p(\phi) := \frac{\omega l(\phi) \cdot p}{\phi}$   $v_g(\phi) := p \cdot \left( \frac{d}{d\phi} \omega l(\phi) \right)$   $\beta_0(\phi) := \frac{\phi}{p}$

$$P(\phi) := \frac{2Lg}{p} \cdot v_g(\phi) \cdot V_a^2 \cdot \left[ \frac{1}{v_{pe}(\phi) \cdot Z_e(\phi)} \cdot (1 + \text{sinc}(k_e(\phi) \cdot Lg)) + \frac{1}{v_{po}(\phi) \cdot Z_e(\phi)} \cdot \left( \frac{\sin(k_e(\phi) \cdot Lg)}{\sin(k_o(\phi) \cdot Lg)} \right) \cdot (1 - \text{sinc}(k_o(\phi) \cdot Lg)) \right] \quad \text{Equation 4.60}$$

$$\text{Electric field in the gap} \quad E_g(\phi, x) := \frac{4 \cdot V_a}{\text{gap}} \cdot \left( \cos(k_e(\phi) \cdot x) \cdot \sin\left(\frac{\phi}{2}\right) + b_{a1}(\phi) \cdot \sin(k_o(\phi) \cdot x) \cdot \cos\left(\frac{\phi}{2}\right) \right) \quad \text{Equation 4.61}$$

$$\text{On the axis the amplitude of the fundamental space harmonic is} \quad E_0(\phi) := E_g(\phi, 0) \cdot \text{sinc}\left(\beta_0(\phi) \cdot \frac{\text{gap}}{2}\right) \cdot \frac{\text{gap}}{p} \quad \text{Equation 4.62}$$

Pierce impedance on the centre line of the structure at height d

$$Z_{p0}(\phi, d) := \frac{E_0(\phi)^2}{2 \cdot \beta_0(\phi)^2 \cdot P(\phi)} \cdot \left[ \frac{\cosh[\beta_0(\phi) \cdot (h_2 - d)]}{\cosh[\beta_0(\phi) \cdot (h_2 - h_1)]} \right]^2 \quad \text{Equation 4.63}$$

$$\text{Plotting ranges} \quad \phi_1 := 0.01 \cdot \pi, 0.02 \cdot \pi \dots \pi \quad x_1 := \frac{-Lg}{2}, -0.8 \cdot \frac{Lg}{2} \dots \frac{Lg}{2}$$

Data from Courtney

Dispersion curve  
(Figure 6)

$$\phi_c := \frac{350}{2 \cdot 134} \cdot \text{deg} \quad f_c := \frac{4}{152}$$

$\begin{pmatrix} 20 \\ 34 \\ 48 \\ 62 \\ 75 \\ 89 \\ 103 \\ 117 \end{pmatrix}$	$\cdot \text{deg}$	$\begin{pmatrix} 35 \\ 56 \\ 75 \\ 91 \\ 107 \\ 121 \\ 134 \\ 144 \end{pmatrix}$
--------------------------------------------------------------------------------	--------------------	----------------------------------------------------------------------------------

Impedance curve  
(Figure 7)

$$Z_{PC} := \begin{pmatrix} 80 \\ 60 \\ 40 \\ 30 \\ 20 \\ 10 \\ 8 \\ 6 \\ 4 \\ 3 \end{pmatrix} \quad f_c := \frac{6}{107}$$

$\begin{pmatrix} 21 \\ 24 \\ 29 \\ 35 \\ 42 \\ 54 \\ 57 \\ 61 \\ 65 \\ 67 \end{pmatrix}$
------------------------------------------------------------------------------------------

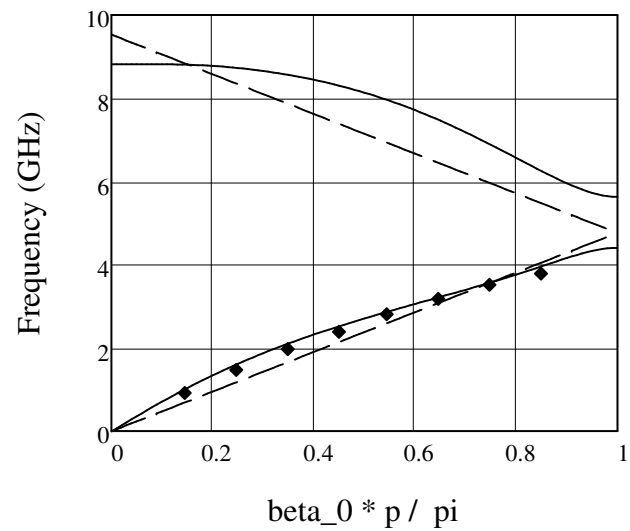


Figure 4.11

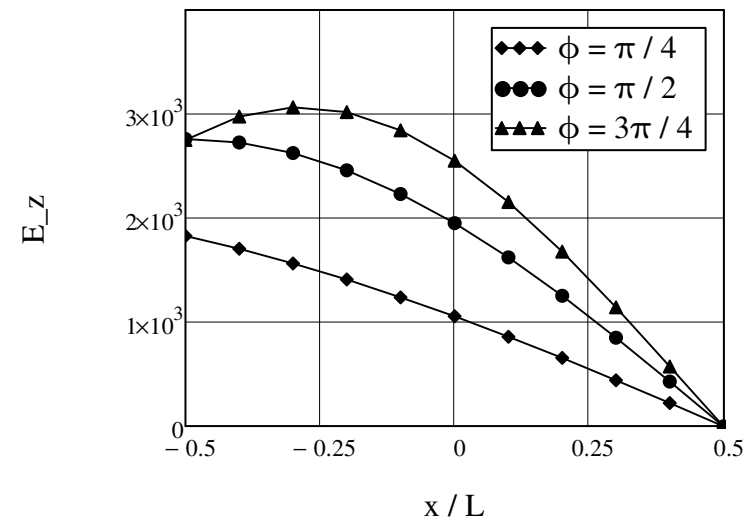
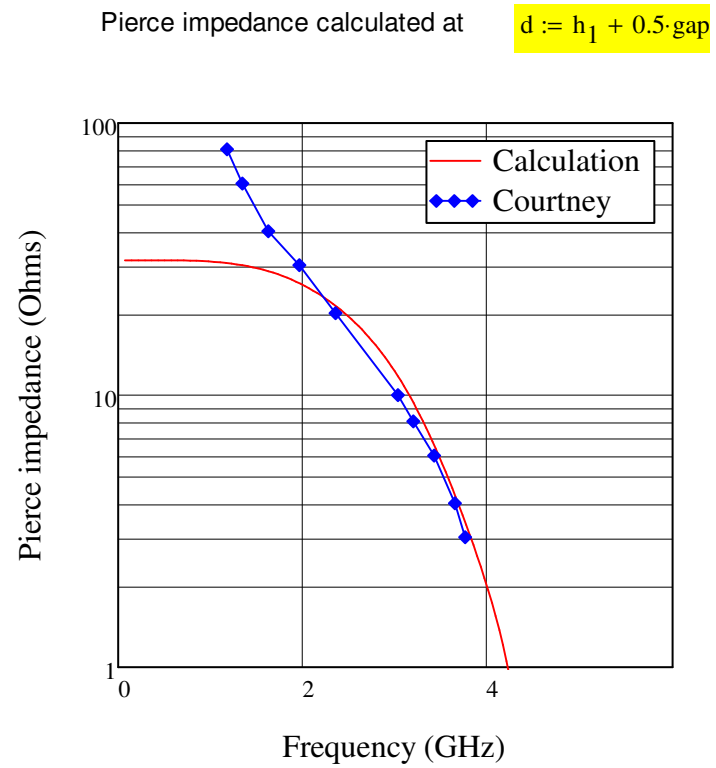


Figure 4.12



Note: It is not quite clear how the impedance in Courtney is defined. Here  $d$  has been adjusted to fit the calculated curve to the experimental data.