## Errata for "An Introduction to Granular Flow"

Last updated: 29 October 2008

- 1. p. xiv, line 4: desribed  $\longrightarrow$  described
- 2. p. xiv, insert after entry for  $k_x$ ,  $k_y$ ,  $k_z$ , the following:  $k_{\rm B}$ : Boltzmann constant
- 3. p. xix, after entry for  $\gamma'$ , insert the following:  $\dot{\gamma}$ : shear rate
- 4. p. xv, after last line, insert the following additional definition for g: magnitude of the relative velocity  $\mathbf{g}$
- 5. p. xvii, after entry for  $\hat{\mathbf{D}}$ , insert the following:  $\hat{\mathbf{D}}'$ :  $\hat{\mathbf{D}} \frac{1}{2}(\hat{\mathbf{\nabla}}\hat{\mathbf{u}})\mathbf{I}$
- 6. p. 38, line above (1.60): article  $\longrightarrow$  particle
- 7. p. 118, line 2 below Table 3.1: bulk density  $\longrightarrow$  density
- 8. p. 141, paragraph 4, line 3:  $\xi^{\lambda-1} \longrightarrow \xi^{\lambda_s-1}$
- 9. p. 162, problem 3.2, last line: Add the following sentence: In the r- $\theta$  plane, the slope of a curve is given by  $\frac{1}{r}\frac{\mathrm{d}r}{\mathrm{d}\theta}$ .
- 10. p. 164, problem 3.9, line 3:  $(3.109) \ll 1 \longrightarrow (3.109)$  is  $\ll 1$
- 11. p. 212: Equation (4.53) should be replaced by

$$(\mathbf{n}_* \cdot \boldsymbol{\sigma})_w = \frac{1}{P'} \oint_{P'} \mathbf{n}_* \cdot \boldsymbol{\sigma} \, \mathrm{d}s$$

- 12. p. 239, paragraph 2 above §5.3, last line: were reported  $\longrightarrow$  were not reported
- 13. p. 277, line above equation (6.57): radial  $\longrightarrow$  incompressible
- 14. p. 301, line 3 from the bottom:  $1.365 \times 10^{-10} \,\mathrm{m} \longrightarrow 2.73 \times 10^{-10} \,\mathrm{m}$
- 15. p. 306, paragraph 2, line 2:  $\langle \psi \rangle / \tau_c \longrightarrow \psi_s / \tau_c$ , where  $\psi_s$  sets the scale of  $\psi$ .
- 16. p. 306, paragraph 2: Replace third sentence, "The other terms ...", by the following:

The other terms, being gradients in the fluxes, are smaller by a factor of s/H. The scaling of the three terms may be derived in the following

manner. Using  $v_s$  as a characteristic grain velocity, we see from (7.65), (7.41) and (7.32) that

$$\chi \sim d_{\rm p}^2 n^2 v_s \psi_s = \frac{6}{\pi} s n \, \frac{\nu}{d_{\rm p}} \, \frac{v_s}{s} \psi_s$$

Here, we have used the identity  $\nu = n\pi d_{\rm p}^3/6$ . Recognizing that  $s \sim d_{\rm p}/\nu$  (see the text below (7.49)) and  $s/v_s = \tau_c$ , we get

$$\chi \sim n \frac{\psi_s}{\tau_c}$$

Similarly, it is easily seen from (7.64) that  $\theta$  scales as

$$\boldsymbol{\theta} \sim d_{\rm p}^3 n^2 v_s \psi_s = \frac{6}{\pi} sn \nu \, \frac{v_s}{s} \psi_s$$

and hence

$$\nabla \cdot \boldsymbol{\theta} \sim \frac{s}{H} n \nu \frac{\psi_s}{\tau_c}$$

In the same manner,

$$\nabla \cdot (n \langle \mathbf{c} \psi \rangle) \sim \frac{1}{H} n v_s \psi_s = \frac{s}{H} n \frac{\psi_s}{\tau_c}$$

Thus, we see that  $\nabla \cdot \boldsymbol{\theta}$  and  $\nabla \cdot (n \langle \mathbf{c} \psi_s \rangle)$  are smaller than  $\chi$  by the factor s/H. Hence, the balance in (7.66) is between the left-hand side and  $\chi(\psi)$ .

17. p. 315, (7.125): 
$$\hat{\mathbf{D}} \longrightarrow \hat{\mathbf{D}}'$$

- 18. p. 316, line 1: with the dimensionless rate of deformation tensor  $\hat{\mathbf{D}} \equiv \hat{\nabla} \hat{\mathbf{v}}$   $\longrightarrow$  with the traceless part of the dimensionless rate of deformation tensor  $\hat{\mathbf{D}}' \equiv \overline{\hat{\nabla}} \hat{\mathbf{v}}$
- 19. p. 317, line 3 below (7.133): that value  $\longrightarrow$  the value

20. p. 318, (7.135): 
$$\hat{\mathbf{D}} \longrightarrow \hat{\mathbf{D}}'$$

21. p. 319, (7.138): 
$$\hat{\mathbf{D}} \longrightarrow \hat{\mathbf{D}}'$$

22. p. 320, (7.142), (7.143), (7.145): 
$$\hat{\mathbf{D}} \longrightarrow \hat{\mathbf{D}}'$$

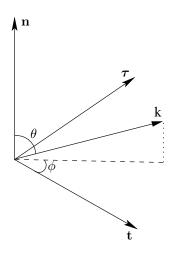
23. p. 320, (7.143) and (7.145): 
$$\nabla \longrightarrow \hat{\nabla}$$

24. p. 323, (7.158): 
$$\mathbf{D} \longrightarrow \mathbf{D}'$$

25. p. 323, line after (7.158): where **D** is the rate of deformation tensor,  $\longrightarrow$  where  $\mathbf{D}' \equiv \mathbf{D} - \frac{1}{3}(\nabla \cdot \mathbf{u}) \mathbf{I}$  is the traceless part of the rate of deformation tensor,

- 26. p. 330, problem 7.3:  $\hat{\mathbf{D}} \longrightarrow \hat{\mathbf{D}}'$
- 27. p. 335, line 4 above (8.11): analogos  $\longrightarrow$  analogous
- 28. p. 345,  $\S 8.2.2$ , line 4: in the hyperbolic  $\longrightarrow$  the hyperbolic
- 29. p. 351,  $\S 8.3.1$ , paragraph 2, line 2: fixed-time  $\longrightarrow$  fixed time
- 30. p. 355, caption of Fig. 8.16, line 1: diameter flowing in a  $\longrightarrow$  diameter in a
- 31. p. 364, paragraph 1, line 2 from the bottom: no matter how large  $\longrightarrow$  for the entire range of
- 32. p. 368, last paragraph, line 6: layers in the x-z plane  $\longrightarrow$  layers parallel to the x-z plane
- 33. p. 368, last line 6: layers are in x-y plane  $\longrightarrow$  layers are parallel to the x-y
- 34. p. 376, 3<sup>rd</sup> line below (9.5):  $r \longrightarrow \mathbf{r}$
- 35. p. 378,  $4^{\text{th}}$  line above (9.13): nearly rough  $\longrightarrow$  nearly perfectly rough
- 36. p. 380, paragraph 2, after (9.27): The body couple, if present . . . from the hydrostatic part.  $\longrightarrow$  If a body couple of finite magnitude is present,  $\overline{\omega}$  is in general not a hydrodynamic variable, but is "enslaved" to  $\tau$ .
- 37. p. 386, paragraph 1: the relations for  $a_1$  and  $b_0$  coincide with the solution obtained for a granular gas of smooth particles, given in (7.133).  $\longrightarrow$  the relation for  $\Phi_K$  differs from the one-term solution of Chapman & Cowling (1964) for a granular gas of smooth particles (given by (7.125), (7.132) and (7.133)) only by the the term proportional to  $a'_1$  in A.
- 38. p. 387, (9.67):  $\hat{\mathbf{D}} \longrightarrow \hat{\mathbf{D}}'$
- 39. p. 387, last line:  $\hat{\mathbf{D}} \longrightarrow \hat{\mathbf{D}}'$
- 40. p. 388, (9.68):  $\hat{\mathbf{D}} \longrightarrow \hat{\mathbf{D}}'$
- 41. p. 389, (9.79):  $\mathbf{D} \longrightarrow \mathbf{D}'$
- 42. p. 389, line below (9.79): where  $\mathbf{D} \equiv -\mathbf{C}$  (see (2.49)) is the rate of deformation tensor defined in the tensile sense.  $\longrightarrow$  where  $\mathbf{D}' \equiv \mathbf{D} \frac{1}{3}(\nabla \cdot \mathbf{u})\mathbf{I}$  is the traceless part of the rate of deformation tensor.
- 43. p. 394, paragraph 2, last line: grains.3  $\longrightarrow$  grains.
- 44. p. 448, line below (G.54):  $[Q]_* \to [Q_*]$

45. p. 450: Fig. H.1 should be modified as below:



- 46. p. 450, §H.2: is evaluated by writing  ${\bf k}$  in terms of  $\longrightarrow$  is evaluated by writing  ${\bf k}$ ,  ${\bf i}$ , and  ${\bf j}$  in terms of
- 47. p. 452, (H.15):  $\mathcal{J} \longrightarrow \det(\mathbf{J})$
- 48. p. 457, line 3 from the bottom: roman  $\longrightarrow$  Roman In the same line: greek  $\longrightarrow$  Greek